An alternative approach to study the heterogeneous dynamics and its consequent dynamic correlation length that emerges in a glass-former.

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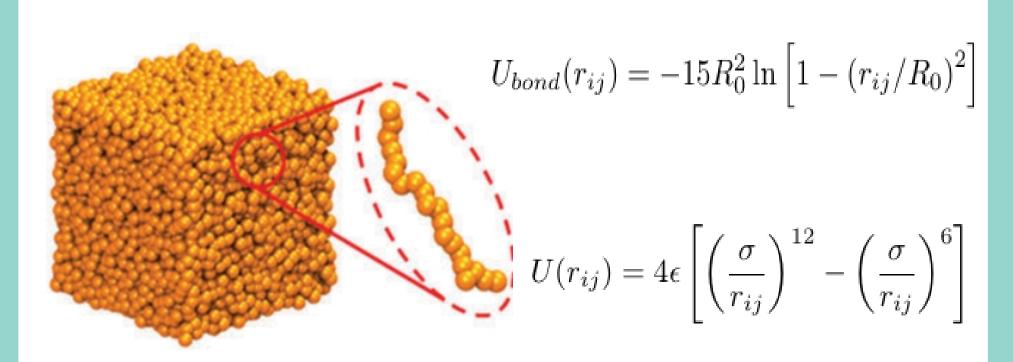
INTRODUCTION

Dynamic slowdown of glass-former liquids, leading to a breakdown of Arrhenius behavior of relaxation and Stokes-Einstein relationship (SER), as the glass transition is approached, is still not fully understood despite decades of study. These are usually associated to the emergence of dynamic heterogeneity (DH). Nevertheless, the physical origin of the DH, and in particular, the question whether they have a structural origin or they are a purely dynamical phenomenon, is still under debate [1].

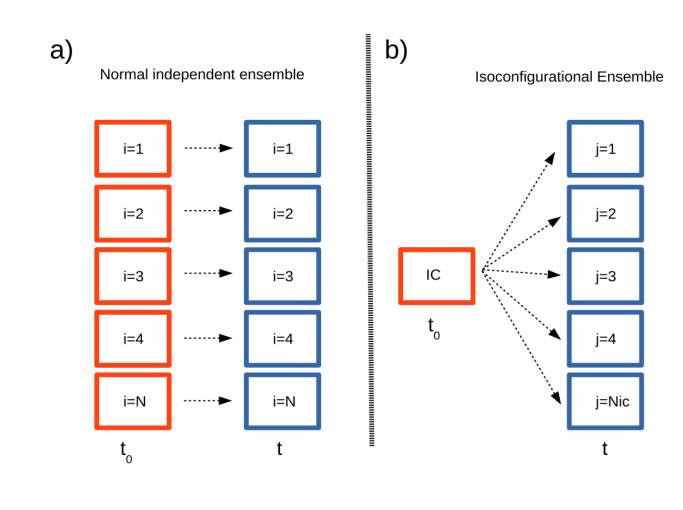
In this work, the relationship between non-Arrhenius behavior, breakdown of SER and DH was analyzed, by means of Molecular Dynamic simulations of a supercooled polymeric glassformer. Our results reinforce the idea that both effects are associated to dynamic heterogeneity related to structural effects. Also, the structural and dynamic correlation lengths that emerge in the system were quantified. The results lead to a consensus of similar scaling of structural and dynamical length scales, reinforcing the idea of the theories of Adam-Gibbs and Random First Order Transition.

METHODS

The attractive Kremer-Grest linear polymer model [2]



- 30 Monomers per chain
- Periodic Boundary Conditions
- LAMMPS Software
- NVE ensemble for production runs
- Isoconfigurational Ensemble (ICE)[3]:



References & Acknowledgments

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0.6

Crossover temperature in the region where the SEB and

the non-Arrhenius behavior are observed

RESULTS We employed the Person's coefficient in the ICE. It quantifies the The N_{corr} from the behavior of dynamics susceptibility was dynamics correlation from the relative displacement of the calculated. particles [4] $\chi_4 = \frac{\beta V}{N^2} [\langle Q(t)^2 \rangle - \langle Q(t) \rangle^2] \longrightarrow N_{corr} = \max_t \{\chi_4(t)\} = \chi_4(t_4)$ $k_{ij}^*(r_{ij}) = \frac{1}{S_i(t^*)S_j(t^*)} \sum_{w=1}^{N_{IC}} \psi_i(w, t^*) \psi_j(w, t^*)$ $\psi_i(w, t^*) = |(\mathbf{r}_i(w, t^*) - \mathbf{r}_i(0))| - \langle \Delta r_i(t^*) \rangle_{IC}$ ∍-⊙ Normal Cluster of particles with high correlation **BREAKDOWN OF STOKES-EINSTEIN** M RELATIONSHIP •--• 0.48 $t^* = \text{maximum}$ in the non-Gaussian parameter $K(r, t^*) \sim \exp{-(r/\xi_K)}$ **△**--**△** 0.54 $t_4 = \text{maximum}$ in the dynamics susceptibility **◆--**0.56 T_x~ change from Arrhenius to super-∘--∘ 0.60 **⋄--•** 0.80 Arrhenius behavior **⊖ ⊕** t* 1.00 **→** 10 --- Arrenenius behaviour ~ 0.75 POINT-TO-SET CORRELATIONS LENGTHS TRUCTURE $q_c(r) - q_0 = \Omega \exp[-((R-a)/\xi_{PtS})^{\zeta}]$ $q_c(R) = \frac{1}{13 \text{ N}} \sum \langle n_i(t_0) n_i(t_0 + \infty) \rangle$ $\Delta \xi_{K}$ S **1.00** q(R) R N PtS with Replica When the IC is considered in the dynamic U Exchange method [5] properties, a linear relationship is obtained for the structure and the dynamic behavior!! without ICEM with ICEM SPECIFIC STRUCTURAL PARAMETERS Bond Orientational Order (Q₆) and Icosahedral fraction (f₁₀) for the Inherent configurations [7] CONCLUSION

We show that the dynamic behavior in this glass-former can be

and dynamical length scales in supercooled liquids.

interpreted through structural behavior. We also highlight that the

alternative way of defining dynamic lengths through the ICE, can be

useful to shed light in the apparent discrepancies between structural